

# Local and Global Well-posedness for the critical Schrödinger-Debye System

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# The Cubic Nonlinear Schrödinger Equation

$$iu_t + \Delta_x u = \epsilon |u|^2 u$$

Here,

- $u : (x, t) \in \mathbb{R}^d \times \mathbb{R} \rightarrow u(x, t) \in \mathbb{C}$ ;
- $\epsilon = 1$  (defocusing) or  $\epsilon = -1$  (focusing).

# Local Well-posedness in $H^s(\mathbb{R}^d)$

Let

$$H^s(\mathbb{R}^d) = \{f \in \mathcal{S}' : \langle \xi \rangle^s \hat{f}(\xi) \in L^2(\mathbb{R}^d)\}, \quad \langle \xi \rangle = \sqrt{1 + |\xi|^2}.$$

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A classical Local Well-Posedness result reads

For all initial data  $u_0 \in H^s(\mathbb{R}^d)$  there exists a unique solution

$$u \in \mathcal{C}([0; T]; H^s(\mathbb{R}^d)), \quad u(0, x) = u_0,$$

where  $T = T(u_0) > 0$ .

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where  $T = T(u_0) > 0$ .

Typically,

$$T = T(\|u_0\|_s) \quad \text{is a decreasing function.}$$

# Local Well-posedness in $H^s(\mathbb{R}^d)$

## The heuristic scaling argument

If  $u$  is a solution, so is  $u_\mu(x, t) = \mu u(\mu x, \mu^2 t)$ . Also,

$$\|u_\mu(t)\|_{\dot{H}^s} = \mu^{s+1-\frac{d}{2}} \|u(\mu^2 t)\|_{\dot{H}^s}.$$

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Indeed, if  $T > 0$  is the lifespan of  $u$ ,

- the lifespan of  $u_\mu$  is  $T_\mu = \frac{T}{\mu^2}$ ;
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If  $s + 1 - \frac{d}{2} < 0$  (ie  $s < s_c$ ),  $T_\mu$  and  $\|u_\mu(0)\|$  both decrease with  $\mu$ .

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**Local Well-Posedness is not expected in  $H^s$  for  $s < s_c$ .**

# Local Well-posedness in $H^s(\mathbb{R}^d)$

This issue has been addressed by many authors during the 80s (Ginibre, Velo, Tsutsumi among others).

The definitive result is the following:

Cazenave and Weissler, 1990

The IVP for the general Schrödinger equation

$$iu_t + \Delta_x u = \pm |u|^2 u$$

is locally well-posed in  $H^s(\mathbb{R}^d)$  for  $s \geq \max\{0; s_c\}$ .

# Global Well-posedness in the energy space $H^1(\mathbb{R}^d)$

The following quantities are formally conserved by the cNLS flow:

$$F(u) = \int |u|^2 \quad E(u) = \frac{1}{2} \int |\nabla u|^2 + \epsilon \frac{1}{4} \int |u|^4.$$

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## Global Well-Posedness: Main idea

Together with LWP, GWP is achieved by controlling  $\|u(t)\|_{H^1}$  (By  $L^2$  conservation, it is enough to control  $\|\nabla u(t)\|_{L^2}$ ).

# Global Well-posedness in the energy space $H^1(\mathbb{R}^d)$

- In the defocusing case this control is immediate:

$$E(u_0) = E(u) = \frac{1}{2} \int |\nabla u|^2 + \frac{1}{4} \int |u|^4 \geq \frac{1}{2} \int |\nabla u|^2.$$



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- In the focusing case, things are not that simple, since

$$E(u) = \frac{1}{2} \int |\nabla u|^2 - \frac{1}{4} \int |u|^4.$$

# Global Well-posedness in the energy space $H^1(\mathbb{R}^d)$ - Focusing case

In dimension  $d = 1$

The Gagliardo Nirenberg inequality

$$\|u\|_{L^4} \leq C \|u\|_{L^2}^{\frac{3}{4}} \|\nabla u\|_{L^2}^{\frac{1}{4}}$$

yields

$$E(u_0) = E(u) \geq \frac{1}{2} \|\nabla u(t)\|_{L^2}^2 - C \|u_0\|_{L^2}^3 \|\nabla u(t)\|_{L^2}$$

which implies  $\|\nabla u(t)\|_{L^2} \leq C$ .

# Global Well-posedness in the energy space $H^1(\mathbb{R}^d)$ - Focusing case

In dimension  $d = 2$

The Gagliardo Nirenberg inequality

$$\|u\|_{L^4} \leq C \|u\|_{L^2}^{\frac{1}{2}} \|\nabla u\|_{L^2}^{\frac{1}{2}}$$

yields

$$E(u_0) = E(u) \geq \frac{1}{2} \|\nabla u(t)\|_{L^2}^2 (1 - C \|u_0\|_{L^2}).$$

which implies  $\|\nabla u(t)\|_{L^2} \leq C$  if  $\|u_0\|_{L^2}$  is small enough.

# Global Well-posedness in the energy space $H^1(\mathbb{R}^d)$ - Focusing case

In dimension  $d = 2, 3$ , the Virial inequality

$$\frac{\partial^2}{\partial t^2} \int |x|^2 |u(x, t)|^2 dx \leq 8dE(u_0)$$

implies blow-up for  $E(u_0) < 0$ .

Note that this can always be achieved: choosing  $u_0 \neq 0$  and taking  $\mu \rightarrow +\infty$ ,

$$E(u_{\mu_0}) = \frac{1}{2}\mu^{2-\frac{d}{2}} \int |\nabla u_0|^2 - \frac{1}{4}\mu^{4-d} \int |u_0|^4 dx \rightarrow -\infty$$

since

$$2 - \frac{d}{2} < 4 - d \Leftrightarrow d < 4.$$

# The Schrödinger-Debye System

$$\begin{cases} iu_t + \frac{1}{2}\Delta u = uv \\ \mu v_t + v = \lambda|u|^2 \end{cases}$$

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Note that for  $\mu = 0$  this system reduces to the cNLS equation.

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- Furthermore, the Schrödinger-Debye system obeys the following pseudo-Hamiltonian structure:

$$\frac{d}{dt}E(t) = 2\lambda\mu \int (v_t)^2 dx,$$

where

$$E(t) = \int (|\nabla u|^2 + 2v|u|^2 - \lambda v^2) dx.$$

# Local Well-Posedness - Previous results

$$\begin{cases} iu_t + \frac{1}{2}\Delta u = uv \\ \mu v_t + v = \lambda|u|^2, \end{cases} \quad u(x, 0) = u_0(x), \quad v(x, 0) = v_0(x).$$

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The second equation can be solved with respect to  $v$ :

$$v(x, t) = e^{-\frac{t}{\mu}} v_0(x) + \frac{\lambda}{\mu} \int_0^t e^{-(t-t')} \mu |u(x, t')|^2 dt'.$$

Hence, the Schrödinger-Debye system can be rearranged into a single integro-differential equation:

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# Local Well-Posedness - Previous results

With this approach:

## B. Bidégaray (1998,2000)

Let  $d = 1, 2, 3$  and  $(u_0, v_0) \in H^s(\mathbb{R}^d) \times H^s(\mathbb{R}^d)$ . Then, there exists  $T > 0$  small enough such that the Schrödinger-Debye system has a unique solution

- $u \in L^\infty([0; T]; H^s(\mathbb{R}^d))$  if  $s > \frac{d}{2}$ ;
- $u \in L^\infty([0; T]; H^1(\mathbb{R}^d))$  if  $s = 1$ ;
- $u \in C([0; T]; L^2(\mathbb{R}^d)) \cap L^{\frac{8}{d}}([0; T]; L^4(\mathbb{R}^d))$  if  $s = 0$ .

# Local Well-Posedness

In the framework of Bourgain spaces introduced by J. Ginibre, Y. Tsutsumi and G. Velo for the Zakharov system, we obtained the following result:

Let  $d = 2, 3$  and  $(u_0, v_0) \in H^s(\mathbb{R}^d) \times H^l(\mathbb{R}^d)$  such that

$$\max\{0, s - 1\} \leq l \leq \min\{2s, s + 1\}.$$

Then, there exists  $T > 0$  small enough such that the Schrödinger-Debye system has a unique solution

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We will shortly describe this method.

# Bourgain spaces

Consider the linear equation

$$iu_t = L(D)u, \quad \widehat{L(D)u}(\xi) = p(\xi)\hat{u}(\xi).$$

(For the Schrödinger equation  $iu_t + \Delta_x u = 0$ ,  $p(\xi) = |\xi|^2$ )

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Taking the Fourier transform in space and time:

$$-\tau\hat{u}(\tau, \xi) = -|\xi|^2\hat{u}(\tau, \xi) : (\tau + |\xi|^2)\hat{u}(\tau, \xi) = 0.$$

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$\hat{u}$  is supported on the paraboloid  $\tau = -|\xi|^2$ .

# Bourgain spaces

Now, we consider a nonlinear perturbation:

$$iu_t = L(D)u + f(u).$$

(For the cNLS equation  $iu_t + \Delta_x u = \pm |u|^2 u$ ,  $f(u) = \pm |u|^2 u$ )

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$\hat{u}$  remains concentrated near the hypersurface  $\tau = -|\xi|^2$ .

We will measure this phenomena using the norm

$$\|u\|_{X^{s,b}}^2 = \| \langle \xi \rangle^s \langle \tau + |\xi|^2 \rangle^b \hat{u}(\tau, \xi) \|_{L^2(\mathbb{R}^{d+1})}^2,$$

$$\langle x \rangle = \sqrt{1 + |x|^2}.$$

# Bourgain spaces

A fundamental result:

## Theorem

Let  $Y = Y(\mathbb{R}^d \times \mathbb{R})$  a Banach space and  $b > \frac{1}{2}$ .

If, for all  $f \in H_x^s$  and for all  $\tau_0 \in \mathbb{R}$ ,

$$\|e^{it\tau_0} e^{-ip(D)t} f\|_Y \leq \|f\|_{H^s}$$

then

$$X^{s,b} \hookrightarrow Y.$$

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## Corollary

$$\text{If } b > \frac{1}{2}, \quad X^{s,b} \hookrightarrow C_t^0 H_x^s.$$

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Hence, (after truncating in time), we will prove the local-posedness of

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Writing this equation in integral form yields ( $S(t) = e^{-ip(D)t}$ )

$$u(t) = S(t)u_0 - i \int_0^t S(t-t')f(u(t'))dt' = S(t)u_0 - iU *_t f(u).$$

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We will deal with

$$u(t) = \psi_T S(t)u_0 - i\psi_T U *_t f(u).$$

Contraction is  $X^{s,b}$ 

$$u(t) = \psi_T S(t)u_0 - i\psi_T U *_t f(u).$$

$$\|u\|_{X^{s,b}} \leq \|\psi_T S(t)u_0\|_{X^{s,b}} + \|\psi_T U *_t f(u)\|_{X^{s,b}}.$$

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- Linear estimates



$$\|\psi_T S(t)u_0\| \leq \|\psi\|_{H_t^b} \|u_0\|_{H_x^s}.$$

- For suitable  $b'$ ,

$$\|\psi_T U *_t f(u)\|_{X^{s,b}} \leq T^{1-b+b'} \|f(u)\|_{X^{s,b'}}.$$



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- Nonlinear estimates

- Estimate  $\|f(u)\|_{X^{s,b'}}$  in terms of  $\|u\|_{X^{s,b}}$

# Back to the Schrödinger-Debye system

We write

$$\begin{cases} iu_t + \frac{1}{2}\Delta u = uv \\ \mu v_t + v = \lambda|u|^2 \end{cases}$$

in integral form:

$$\begin{cases} u(t) = S(t)u_0 - i \int_0^t S(t-t')uv(t')dt' \\ v(t) = e^{-\frac{t}{\mu}}v_0 + \frac{\lambda}{\mu} \int_0^t e^{\frac{-(t-t')}{\mu}}|u(t')|^2 dt' \end{cases}$$

and define Bourgain spaces adapted to  $u$  and  $v$ :

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and define Bourgain spaces adapted to  $u$  and  $v$ :

$$\|u\|_{X^{s,b}} = \left\| \langle \xi \rangle^s \langle \tau + \frac{1}{2}|\xi|^2 \rangle^b \hat{u}(\xi, \tau) \right\|_{L^2};$$

$$\|v\|_{H^{l,c}} = \left\| \langle \xi \rangle^l \langle \tau \rangle^c \hat{v}(\xi, \tau) \right\|_{L^2}.$$

# Back to the Schrödinger-Debye system

As explained, one only needs to estimate the nonlinear terms

- $\|f_1(u, v)\|_{X^{s,b'}} = \|uv\|_{X^{s,b'}}$
- $\|f_1(u, v)\|_{H^{l,c'}} = \| |u|^2 \|_{X^{l,c'}}$

by the norms  $\|u\|_{X^{s,b}}$  and  $\|v\|_{H^{l,c}}$ .

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by the norms  $\|u\|_{X^{s,b}}$  and  $\|v\|_{H^{l,c}}$ .

We established the following (bilinear) estimates:

$$\|uv\|_{X^{s,-\frac{1}{2}+}} \leq C \|u\|_{X^{s,\frac{1}{2}+}} \|v\|_{H^{l,\frac{1}{2}+}}, \quad s \geq 0, l \geq \max\{0, s-1\};$$

$$\|u\bar{w}\|_{H^{l,-\frac{1}{2}+}} \leq C \|u\|_{X^{s,\frac{1}{2}+}} \|w\|_{X^{s,\frac{1}{2}+}}, \quad s \geq 0, l \leq \min\{2s, s-1\};$$

# Global well-posedness of the D-S system

A. Corcho, J.D. Silva & FO (2011)

Let  $(u_0, v_0) \in H^1(\mathbb{R}^2) \times L^2(\mathbb{R})$ . Then, for all  $T > 0$ , there exists a unique solution

$$(u, v) \in C([0; T], H^1(\mathbb{R}^2) \times L^2(\mathbb{R})).$$

to the I.V.P. associated to the Debye-Schrödinger system.

This theorem remains valid in both focusing and defocusing cases  $\lambda = \pm 1$ .

# Global well-posedness of the D-S system

Idea of the proof:

Compute an a priori bound for the quantity

$$f(t) = \|\nabla u(t)\|_{L^2}^2 + \|v(t)\|_{L^2}^2.$$

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and  $\|\nabla u\|_{L^2}^2$  in terms of the pseudo-energy  $E(t)$ :

$$\|\nabla u\|_{L^2}^2 = E(t) - \int (2v|u|^2 - \lambda v^2) dx.$$

# Global well-posedness of the D-S system

Example of estimate:

$$\begin{aligned}\|v\|_2 &\leq \|v_0\|_2 + \frac{1}{\mu} \int_0^t e^{-\frac{(t-t')}{\mu}} \|u(t')\|_{L^4}^2 dt' \\ &\leq \|v_0\|_2 + \frac{C^2}{\mu} \int_0^t \|u(t')\|_{L^2} \|\nabla u(t')\|_{L^2} dt' \\ &\leq \|v_0\|_2 + \frac{C^2 \|u_0\|_{L^2}}{\mu} \int_0^t \|\nabla u(t')\|_{L^2} dt'\end{aligned}$$

# Global well-posedness of the D-S system

By squaring,

$$\|v\|_{L^2}^2 \leq 2\|v_0\|_2^2 + \frac{C^4 \|u_0\|_{L^2}^2}{\mu^2} \left( \int_0^t \|\nabla u(t')\|_{L^2} dt' \right)^2$$

and

$$\|v\|_{L^2}^2 \leq 2\|v_0\|_2^2 + \frac{C^4 \|u_0\|_{L^2}^2}{\mu^2} t \int_0^t \|\nabla u(t')\|_{L^2}^2 dt'$$

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Finally,

$$\|v\|_{L^2}^2 \leq 2\|v_0\|_2^2 + \frac{C^4 \|u_0\|_{L^2}^2}{\mu^2} t \int_0^t f(t') dt'.$$

# Global well-posedness of the D-S system

One can show that

For all  $t \leq T_\mu = T_\mu(\|u_0\|_2)$ ,

$$f(t) \leq \alpha_0 + \alpha_1 \int_0^t f(t') dt',$$

where  $\alpha_0$  and  $\alpha_1$  depend exclusively on the initial data.