

# On a discrete Boltzmann equation

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(joint work with A.J.Soares, U. Minho)

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- 1 Discrete Boltzmann Equations
- 2 A discrete model for a chemically active gas
- 3 Local solutions
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# Boltzmann 1870

Description of the dynamics of a rarefied ideal gas.  $f(\vec{x}, \vec{v}, t)$ :  
distribution function of particles in the phase space  
(position, momentum) at time  $t$ .

$$\frac{Df}{Dt} = \frac{\partial f}{\partial t}(\vec{x}, \vec{v}, t) + \vec{v} \cdot \vec{\nabla}_{\vec{x}} f(\vec{x}, \vec{v}, t) = \langle f, f \rangle .$$

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Collision operator:

$$\langle f, f \rangle = \int \int_{\Omega, \vec{v}_2} (f^* f_2^* - f f_2) |\vec{v} - \vec{v}_2| \sigma(\Omega) d\Omega d\vec{v}_2,$$

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where

$$f = f(\vec{x}, \vec{v}, t), \quad f_2 = f(\vec{x}, \vec{v}_2, t), \quad f^* = f(\vec{x}, \vec{v}^*, t), \quad f_2^* = f(\vec{x}, \vec{v}_2^*, t),$$

$\vec{v}^*, \vec{v}_2^*$  functions of  $\vec{v}, \vec{v}_2, \Omega$ .

Statistical Entropy (Boltzmann, Gibbs 1872).

Milestone: Boltzmann  $\mathcal{H}$  function,

$$\mathcal{H}(\vec{x}, t) = \int_{\mathcal{V}} f(\vec{x}, \vec{v}, t) \log(f(\vec{x}, \vec{v}, t)) dv$$

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$\mathcal{H}$ -Theorem

$$\frac{\partial}{\partial t} \mathcal{H}(\vec{x}, t) \leq 0.$$

# Discrete Boltzmann Models: the Broadwell Model (1964)

We will allow particles to travel at a finite number of preselected velocities only.

6-velocity Broadwell Model:

$$v_1 = (c, 0, 0), \quad v_2 = (0, c, 0), \quad v_3 = (0, 0, c)$$

and  $v_{j+3} = -v_j$  for  $j = 1, 2, 3$ .



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$N_i(\vec{x}, t)$ : Number density of particles travelling with speed  $v_i$ .

# Evolution system for $N = (N_1, \dots, N_6)$

Admissible (inelastic) collisions:

Conservation of kinetic energy and momentum.

Particles with velocity  $v_i$  collide with particles with velocity  $-v_i = v_{i+3}$  with 3 possible outcomes (each with probability  $p = 1/3$ .)

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We will focus on the one-dimensional evolution. The Broadwell system becomes:

$$\left\{ \begin{array}{l} \frac{\partial N_1}{\partial t} + c \frac{\partial N_1}{\partial x} = \frac{4cS}{3}(N_2^2 - N_1 N_3) \\ \frac{\partial N_2}{\partial t} = \frac{2cS}{3}(N_1 N_3 - N_2^2) \\ \frac{\partial N_3}{\partial t} - c \frac{\partial N_3}{\partial x} = \frac{4cS}{3}(N_2^2 - N_1 N_3) \end{array} \right.$$

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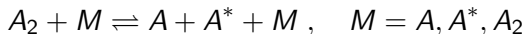
Kawashima (Nonlinear analysis-TMA, 1990):

Existence of strong global solutions for this system.

# A model for a chemically active gas

1997: R. Monaco, M. Pandolfi Bianchi and A.J. Soares.

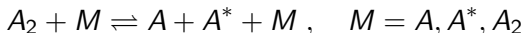
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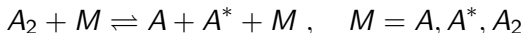
**Kinetic equations:**  $(v_1, v_2, v_3) = (c, 0, -c)$

$$\left\{ \begin{array}{l} \frac{\partial}{\partial t} N_i^A(x, t) + v_i \frac{\partial}{\partial x} N_i^A(x, t) = F_i^A(N(x, t)) \quad (i \in \{1; 2; 3\}), \\ \frac{\partial}{\partial t} N_i^{A_2}(x, t) + \frac{v_i}{2} \frac{\partial}{\partial x} N_i^{A_2}(x, t) = F_i^{A_2}(N(x, t)) \quad (i \in \{1; 2; 3\}), \\ \frac{\partial}{\partial t} N_i^{A^*}(x, t) + v_i \frac{\partial}{\partial x} N_i^{A^*}(x, t) = F_i^{A^*}(N(x, t)) \quad (i \in \{1; 3\}), \end{array} \right.$$

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$$N = \left( N_1^A, N_2^A, N_3^A, N_1^{A^*}, N_3^{A^*}, N_1^{A_2}, N_2^{A_2}, N_3^{A_2} \right).$$

Structure of the collision terms:

$$F_i^M(N) = \left( P_{i,M}^{(1)}(N) + P_{i,M}^{(2)}(N) \right) - N_i^M \left( Q_{i,M}^{(1)}(N) + Q_{i,M}^{(2)}(N) \right)$$

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- $P_{i,M}^{(1)}(N)$  and  $P_{i,M}^{(2)}(N)$ : polynomials of degree 2 and 3, representing the creation of particles  $M$  with velocity  $v_i$  due to inert or reactive collisions.

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- $Q_{i,M}^{(1)}(N)$  and  $Q_{i,M}^{(2)}(N)$ : polynomials of degree 1 and 2, representing the disappearance of particles  $M$  with velocity  $v_i$  due to inert or reactive collisions.



We will study the mixed problem in the half-space:  
 $(x, t) \in [0; +\infty[ \times [0; +\infty[.$

- Initial conditions

$$N_i^M(x, 0) = N_{i_0}^M(x), \quad x \geq 0$$

- Boundary conditions

$$\begin{pmatrix} N_1^A(0, t) \\ N_1^{A^*}(0, t) \\ N_1^{A_2}(0, t) \end{pmatrix} = \begin{pmatrix} \beta_A^A & \beta_{A^*}^A & \beta_{A_2}^A \\ \beta_A^{A^*} & \beta_{A^*}^{A^*} & \beta_{A_2}^{A^*} \\ \beta_A^{A_2} & \beta_{A^*}^{A_2} & \beta_{A_2}^{A_2} \end{pmatrix} \begin{pmatrix} N_3^A(0, t) \\ N_3^{A^*}(0, t) \\ N_3^{A_2}(0, t) \end{pmatrix}$$

$$\sum_{M'} \beta_{M'}^M \leq 1, \quad \beta_M^A + \beta_M^{A^*} + \frac{1}{2}\beta_M^{A_2} \leq \delta_M,$$

where  $\delta_M = 1$  if  $M = A, A^*$  and  $\delta_{A_2} = \frac{1}{2}$ .

We set

$$\mathcal{B}^1(X) = C^1(X) \cap W^{1,\infty}(X) = \left\{ f \in C^1(X) / f, \frac{\partial f}{\partial x}, \frac{\partial f}{\partial t} \in L^\infty(X) \right\}.$$

### Local solutions

Let  $N_o \in \mathcal{B}_+^1([0; +\infty[)$ . Then there exists a unique solution

$$N \in \mathcal{B}_+^1([0; +\infty[ \times [0; T_o])$$

for the mixed-problem, where the life-span  $T_o > 0$  depends exclusively on

$$E_o = \max_{i,M} \sup_{x \in \mathbb{R}^+} N_{i_o}^M(x).$$

## sketch of the proof

We will use a Banach fixed-point technique.

$$\frac{\partial}{\partial t} N_i^M + v_i^M \frac{\partial}{\partial x} N_i^M = F_i^M(N).$$

We sum  $\lambda N_i^M$ :

$$\frac{\partial}{\partial t} N_i^M + v_i^M \frac{\partial}{\partial x} N_i^M + \lambda N_i^M = F_i^M(N) + \lambda N_i^M := F_i^{M,\lambda}(N) \quad (1)$$

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We fix  $(x, t)$  and integrate (1) along the  $v_i^M$ -characteristic through  $(x, t)$ .

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1. If  $x - v_i^M(t) \geq 0$ :

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2. If  $x - v_i^M(t) < 0$ :

$$N_i^M(x, t) = e^{-\lambda(t-t_i^M)} N_i^M(0, t_i^M) + \int_{t_i^M}^t e^{-\lambda(t-\tau)} F_i^{M,\lambda}(N)(x - v_i^M(t - \tau), \tau) d\tau.$$

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We use the boundary condition:



$$N_i^M(x, t) = e^{-\lambda(t-t_i^M)} \sum_{M'} \beta_{M'}^M N_3^{M'}(0, t_i^M) + \int_{t_i^M}^t e^{-\lambda(t-\tau)} F_i^{M,\lambda}(N)(x - v_i^M(t - \tau), \tau) d\tau.$$

We integrate once again along the characteristics:

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$$N_i^M(x, t) = e^{-\lambda t} \sum_{M'} \beta_{M'}^M e^{-\lambda t_i^M} N_{3_0}^{M'}(-v_3^{M'} t_i^M) + \\ + e^{-\lambda t} \sum_{M'} \beta_{M'}^M \int_0^{t_i^M} e^{-\lambda(t_i^M-\tau)} F_3^{M',\lambda}(N)(x - v_3^{M'}(t-\tau), \tau) d\tau \\ + \int_{t_i^M}^t e^{-\lambda(t-\tau)} F_i^{M,\lambda}(N)(x - v_i^M(t-\tau), \tau) d\tau := \text{expr}_2(N).$$

We set:

- $\Omega_T = [0; +\infty[ \times [0; T]$   
 $\Omega_{M,i,T}^+ = \{(x, t) \in \Omega_T ; x \geq v_i^M t\}$   
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- $X^1(T) = \left\{ N = (N_i^M)_{i,M} ; N, \frac{\partial N_i^M}{\partial x} \in C^0 \cap L^\infty(\Omega_T) \right\}$
- $S(T, E, G) \subset X^1(T)$  the close and convex set of functions  $N$  such that  
 $\forall i, M$ 
  - $N_i^M(x, 0) = N_{i_0}^M(x) , x \in [0; +\infty[$
  - $0 \leq N_i^M(x, t) \leq E , \left| \frac{\partial N_i^M}{\partial x}(x, t) \right| \leq G , (x, t) \in \Omega_T.$

We now set

- $\psi : S(T, E, G) \rightarrow S(T, E, G)$   
 $N_i^M(x, t) \rightarrow \text{expr}_1(N)$  if  $(x, t) \in \Omega_{M,i,T}^+$   
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We can prove that for adequate values of  $S, T, G, \lambda$ ,  $\psi$  is a contraction of  $S(T, E, G)$ .

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The time regularity is easy to get.



## Global solutions (2005)

Let  $N_o \in \mathcal{B}_+^1([0; \infty[) \cap L^1([0; +\infty[)$ .

Then there exists  $\epsilon, \epsilon' > 0$  such that if

$$m_o = \int_0^{+\infty} \left( \sum_{i=1}^3 (N_{i_o}^A + 2N_{i_o}^{A_2}) + \sum_{i=1,3} N_{i_o}^{A^*} \right) (x) dx < \epsilon$$

and

$$E_o = \max_{i,M} \sup_{x \in \mathbb{R}^+} N_{i_o}^M(x) < \epsilon',$$

the mixed problem has a unique solution

$$N \in \mathcal{B}_+^1([0; \infty[ \times [0; \infty[.$$

**Proof:** We set  $E(T_o) = \max_{i,M} \sup_{x \in \mathbb{R}^+, t \in [0; T_o]} N_i^M(x, t)$ .

We show an a priori estimate of the type  $\forall t \in [0; T], E(t) \leq M$ .

## Conservation laws:



$$\frac{\partial}{\partial t} \left( \sum_{i=1}^3 (N_i^A + N_i^{A_2}) \right) + \frac{\partial}{\partial x} \left( \sum_{i=1}^3 v_i (N_i^A + \frac{1}{2} N_i^{A_2}) \right) = 0$$



$$\frac{\partial}{\partial t} \left( \sum_{i=1}^3 N_i^{A_2} + \sum_{i=1,3} N_i^{A^*} \right) + \frac{\partial}{\partial x} \left( \sum_{i=1}^3 \frac{1}{2} v_i N_i^{A_2} + \sum_{i=1,3} v_i N_i^{A^*} \right) = 0$$



$$\frac{\partial}{\partial t} \left( \sum_{i=1}^3 v_i (N_i^A + N_i^{A_2}) + \sum_{i=1,3} v_i N_i^{A^*} \right) + \frac{\partial}{\partial x} \left( \sum_{i=1}^3 v_i^2 (N_i^A + \frac{1}{2} N_i^{A_2}) + \sum_{i=1,3} v_i^2 N_i^{A^*} \right) = 0.$$

From the conservation laws, we build two exact forms:

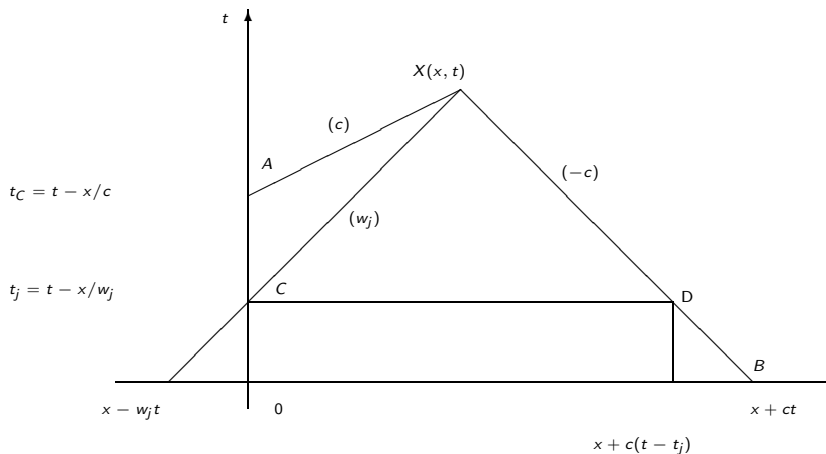
$$w = - \left( \sum_{i=1}^3 (N_i^A + 2N_i^{A_2}) + \sum_{i=1,3} N_i^{A^*} \right) dx +$$

$$\left( \sum_{i=1}^3 v_i (N_i^A + N_i^{A_2}) + \sum_{i=1,3} v_i N_i^{A^*} \right) dt.$$

$$w' = - \left( \sum_{i=1}^3 ((w_j - v_i) N_i^A + (2w_j - v_i) N_i^{A_2}) + \sum_{i=1,3} (w_j - v_i) N_i^{A^*} \right) dx$$

$$+ \left( \sum_{i=1}^3 v_i ((w_j - v_i) N_i^A + (w_j - \frac{1}{2} v_i) N_i^{A_2}) + \sum_{i=1,3} v_i (w_j - v_i) N_i^{A^*} \right) dt.$$

We integrate the 1-forms  $w$  and  $w'$  in cycles:



, Eventually, we will get the inequality

$$E(t) \leq CE_o + C'm_o(E(t) + E(t)^2).$$