

# NONTRIVIAL GROUND STATES FOR COOPERATIVE CUBIC SCHRÖDINGER SYSTEMS

---

Fiipe Oliveira

Tomar, 19 May, 2016

CMA, Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa

**FCT**

Fundação para a Ciência e a Tecnologia  
MINISTÉRIO DA CIÊNCIA, INOVAÇÃO E DO TERCEIRO SECTOR

UID/MAT/00297/2013

**FCT** FACULDADE DE  
CIÊNCIAS E TECNOLOGIA  
UNIVERSIDADE NOVA DE LISBOA

We will consider the following coupling of Schrödinger systems:

$$\begin{cases} i\partial_t \phi_j - \Delta \phi_j = \mu_j |\phi_j|^{2q-2} \phi_j + \phi_j |\phi_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |\phi_k|^q \\ \phi_j \in H_0^1(\Omega, \mathbb{C}), \quad j = 1, \dots, n \end{cases}$$

We will consider the following coupling of Schrödinger systems:

$$\begin{cases} i\partial_t \phi_j - \Delta \phi_j = \mu_j |\phi_j|^{2q-2} \phi_j + \phi_j |\phi_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |\phi_k|^q \\ \phi_j \in H_0^1(\Omega, \mathbb{C}), \quad j = 1, \dots, n \end{cases}$$

One important feature of these systems is the existence of so-called **bound-state solutions**, of the form

$$\phi_j(t, x) = e^{-i\lambda_j t} u_j(x).$$

We will consider the following coupling of Schrödinger systems:

$$\begin{cases} i\partial_t \phi_j - \Delta \phi_j = \mu_j |\phi_j|^{2q-2} \phi_j + \phi_j |\phi_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |\phi_k|^q \\ \phi_j \in H_0^1(\Omega, \mathbb{C}), \quad j = 1, \dots, n \end{cases}$$

One important feature of these systems is the existence of so-called **bound-state solutions**, of the form

$$\phi_j(t, x) = e^{-i\lambda_j t} u_j(x).$$

$$\begin{cases} -\Delta u_j + \lambda_j u_j = \mu_j |u_j|^{2q-2} u_j + u_j |u_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |u_k|^q & \text{in } \Omega \\ u_1, \dots, u_n \in H_0^1(\Omega) \end{cases}$$

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

There is an associated energy:

$$\begin{aligned} E(u, v) &= \int \left( \lambda_1 |u|^2 + \lambda_2 |v|^2 + |\nabla u|^2 + |\nabla v|^2 \right) \\ &\quad - \frac{1}{2q} \int \left( |u|^{2q} + |v|^{2q} + 2b |uv|^q \right). \end{aligned}$$

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

There is an associated energy:

$$\begin{aligned} E(u, v) &= \int \left( \lambda_1 |u|^2 + \lambda_2 |v|^2 + |\nabla u|^2 + |\nabla v|^2 \right) \\ &\quad - \frac{1}{2q} \int \left( |u|^{2q} + |v|^{2q} + 2b |uv|^q \right). \end{aligned}$$

**Ground-State Solution** Among all the bound-states, the ones which minimize the Energy.

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

There is an associated energy:

$$\begin{aligned} E(u, v) &= \int \left( \lambda_1 |u|^2 + \lambda_2 |v|^2 + |\nabla u|^2 + |\nabla v|^2 \right) \\ &\quad - \frac{1}{2q} \int \left( |u|^{2q} + |v|^{2q} + 2b |uv|^q \right). \end{aligned}$$

**Ground-State Solution** Among all the bound-states, the ones which minimize the Energy.

They are said **Fully Nontrivial** if all components are non-null.

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

There is an associated energy:

$$\begin{aligned} E(u, v) &= \int \left( \lambda_1 |u|^2 + \lambda_2 |v|^2 + |\nabla u|^2 + |\nabla v|^2 \right) \\ &\quad - \frac{1}{2q} \int \left( |u|^{2q} + |v|^{2q} + 2b |uv|^q \right). \end{aligned}$$

**Ground-State Solution** Among all the bound-states, the ones which minimize the Energy.

They are said **Fully Nontrivial** if all components are non-null.

Otherwise, they are said **Semitrivial**.

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

**Theorem (FO, 2015)**

- If  $1 < q < 2$ , then for all  $b > 0$  the system admits Fully Nontrivial Ground States;
- If  $q \geq 2$  and  $\beta \geq \frac{2^q-1}{2} \omega^{1+\frac{q}{2}} - \frac{1}{2} \omega^{-\frac{q}{2}}$ , where  $\omega = \lambda_2/\lambda_1$ , then the system admits Fully Nontrivial Ground States.

$n = 2$  equations:

$$\begin{cases} -\Delta u + \lambda_1 u = \mu_1 |u|^{2q-2} u + \beta u |u|^{q-2} |v|^q & \text{in } \Omega \\ -\Delta v + \lambda_2 v = \mu_2 |v|^{2q-2} v + \beta v |v|^{q-2} |u|^q & \text{in } \Omega \\ u_1, u_2 \in H_0^1(\Omega) \end{cases}$$

### Theorem (FO, 2015)

- If  $1 < q < 2$ , then for all  $b > 0$  the system admits Fully Nontrivial Ground States;
- If  $q \geq 2$  and  $\beta \geq \frac{2^q-1}{2} \omega^{1+\frac{q}{2}} - \frac{1}{2} \omega^{-\frac{q}{2}}$ , where  $\omega = \lambda_2/\lambda_1$ , then the system admits Fully Nontrivial Ground States.

Referee Report: The author is invited to have a look at the following publication containing all results which the author proved in his submission: Mandel, R. : Minimal energy solutions for cooperative nonlinear Schrödinger systems, *Nonlinear Differential Equations and Applications NoDEA*, <http://dx.doi.org/10.1007/s00030-014-0281-2>.

The general system

$$\begin{cases} -\Delta u_j + \lambda_j u_j = \mu_j |u_j|^{2q-2} u_j + u_j |u_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |u_k|^q & \text{in } \Omega \\ u_1, \dots, u_n \in H_0^1(\Omega) \end{cases}$$

The general system

$$\begin{cases} -\Delta u_j + \lambda_j u_j = \mu_j |u_j|^{2q-2} u_j + u_j |u_j|^{q-2} \sum_{k \neq j} \beta_{j,k} |u_k|^q & \text{in } \Omega \\ u_1, \dots, u_n \in H_0^1(\Omega) \end{cases}$$

### **Theorem**

**(Hugo Tavares, F.O., Advanced Nonlinear Studies, 2016)**

Let  $\lambda_i, \mu_i, \beta_{j,k} > 0$ . Then, for  $1 < q < 2$ , the system admits a Fully Nontrivial Ground State.

For  $q = 2$  and  $k \geq 3$  equations, the situation is more complex.

$$-\Delta u_i + \lambda_i u_i = \mu_i u_i^3 + \beta u_i \sum_{j \neq i} u_j^2 \text{ in } \mathbb{R}^N$$

### First Existence Results of Nontrivial Groundstates

Theorem (Z. Liu, Z.-Q. Wang, 2010)

$$\lambda_1 = \dots = \lambda_k \quad \beta > k(k-1) \max\{\mu_i\} - \frac{d-1}{d} \sum_j \mu_j$$

*Then all groundstates are nontrivial.*

For  $q = 2$  and  $k \geq 3$  equations, the situation is more complex.

$$-\Delta u_i + \lambda_i u_i = \mu_i u_i^3 + \beta u_i \sum_{j \neq i} u_j^2 \text{ in } \mathbb{R}^N$$

### First Existence Results of Nontrivial Groundstates

Theorem (Z. Liu, Z.-Q. Wang, 2010)

$$\lambda_1 = \dots = \lambda_k \quad \beta > k(k-1) \max\{\mu_i\} - \frac{d-1}{d} \sum_j \mu_j$$

*Then all groundstates are nontrivial.*

Later improved to:

Theorem (H. Liu, Z. Liu, J. Chang 2015)

$$\lambda_1 = \dots = \lambda_k \quad \beta > \max\{\mu_j\}.$$

*Then all groundstates are nontrivial.*

## Nontrivial groundstates for different $\lambda_i$ 's

Theorem (Correia, Tavares, F.O., 2015)

Let  $k \geq 3$ ,  $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n$ . Then there exists

$$\alpha = \alpha(\lambda_1/\lambda_2, n, N) < 1$$

such that, if

$$\lambda_n \leq \alpha \lambda_2,$$

then there exists a constant  $B = B(\lambda_i, \mu_i) > 0$ , such that, for  $\beta > B$ , all ground states are nontrivial.

Idea:  $\lambda_1$  and  $\lambda_2$  can be arbitrarily far, but all other  $\lambda_i$  need to be close to  $\max\{\lambda_1, \lambda_2\}$ .

## Nontrivial groundstates for different $\lambda_i$ 's

Theorem (Correia, Tavares, F.O., 2015)

Let  $k \geq 3$ ,  $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_n$ . Then there exists

$$\alpha = \alpha(\lambda_1/\lambda_2, n, N) < 1$$

such that, if

$$\lambda_n \leq \alpha \lambda_2,$$

then there exists a constant  $B = B(\lambda_i, \mu_i) > 0$ , such that, for  $\beta > B$ , all ground states are nontrivial.

Idea:  $\lambda_1$  and  $\lambda_2$  can be arbitrarily far, but all other  $\lambda_i$  need to be close to  $\max\{\lambda_1, \lambda_2\}$ .

Is this result optimal?

In a qualitative way, yes!

Theorem (Correia, Tavares, FO 2015)

Let  $k \geq 3$ ,  $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_k$ .

There exists a constant  $\Lambda = \Lambda(\lambda_1/\lambda_2)$  such that, if

$$\lambda_2 \Lambda \leq \lambda_i \text{ for some } i \geq 3, \text{ and } \beta > \max\{\mu_1, \dots, \mu_k\},$$

then every groundstate solution is semitrivial.

In a qualitative way, yes!

Theorem (Correia, Tavares, FO 2015)

Let  $k \geq 3$ ,  $0 < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_k$ .

There exists a constant  $\Lambda = \Lambda(\lambda_1/\lambda_2)$  such that, if

$$\lambda_2 \Lambda \leq \lambda_i \text{ for some } i \geq 3, \text{ and } \beta > \max\{\mu_1, \dots, \mu_k\},$$

then every groundstate solution is semitrivial.

Dear authors,

Your manuscript *Semitrivial vs. fully nontrivial ground states in cooperative cubic Schrödinger systems with  $d \geq 3$  equations* has been evaluated by two referees. Both evaluators have been very positive and I am glad to announce that your manuscript is accepted.(...)

Cedric Villani

Editor, Journal of Functional Analysis